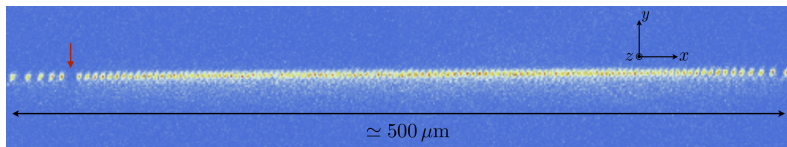


# Experimental Progress on Quantum Computing with Atomic Qubits

## Xiang Zhang



中國人民大學  
RENMIN UNIVERSITY OF CHINA

January 26, 2018

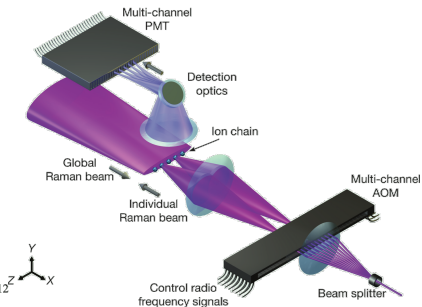
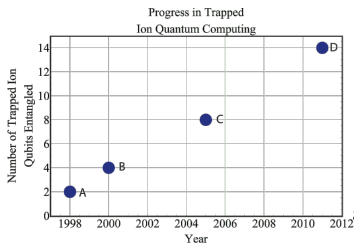
# Outline

- 1 Trapped-ion System
- 2 Experimental Violation of Quantum Contextuality
- 3 Symmetry Operations with an Embedding Quantum Simulator
- 4 Quantum Simulation of Quantum Field Theory
- 5 Conclusion

# Trapped Ions: Promising Architecture<sup>1</sup>

## Scalable and universal trapped-ions quantum computer

- Long coherence time: up to seconds or even hours
- Perfect quantum operation: fidelities of gate and measurement  $> 99\%$
- Local scalability: shuttling or addressing  $> 10$  ions in one trap
- Quantum networks: remotely entangled ion chains through photons



<sup>1</sup>S. Debnath, et. al., Nature, 536, 63–66 (2016)

# Ion Trap

- Captures and confines ions in a vacuum system
- Precision measurement: most accurate atomic clock, gyroscope
- Ionization and control: mass spectrometer, vacuum pump/gauge
- Penning trap: an axial magnetic ring and two endcaps
- Paul trap: four RF electrodes and two DC needles

## Perfect pure quantum system

- Isolated system with ultra-high vacuum ( $< 10^{-11}$  torr)
- Atomic levels and well designed harmonic trapping potential
- Universal rich set of quantum operations

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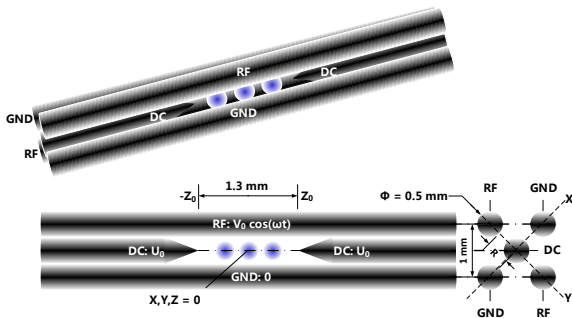
## Perfect pure quantum system

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# 4-rod Trap

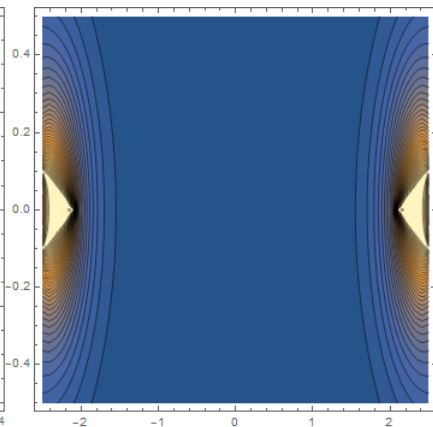
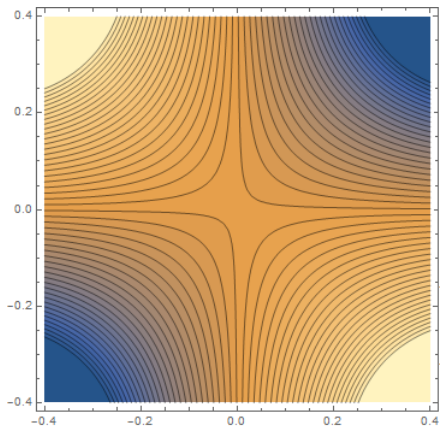
## Trap design

- 4 rods: parabolic pseudopotential formed by rotating RF field
- 2 needles: static Coulomb potential
- 2 extra electrodes: compensate background electric field
- Ions arranged into a linear string



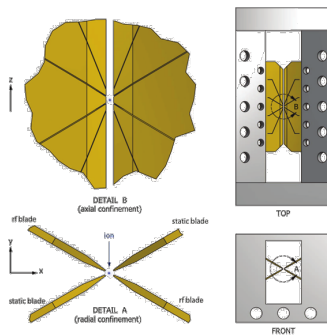
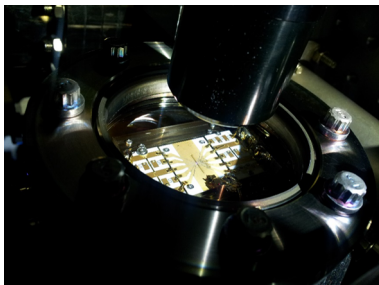
# 4-rod Trap

Calculating trap's X-Y and Z potential with BEM method



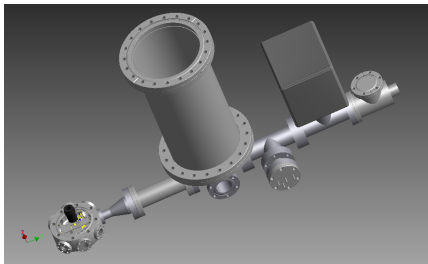
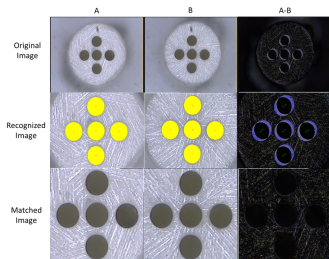
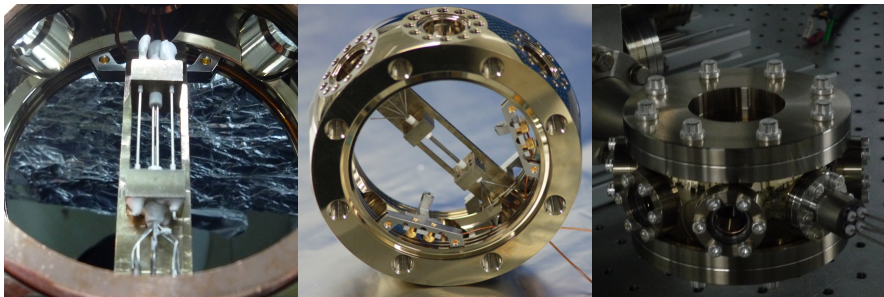
# New Traps

## New trap designs

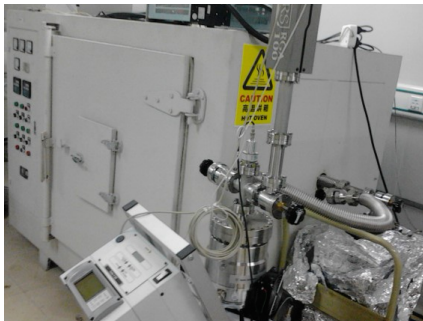
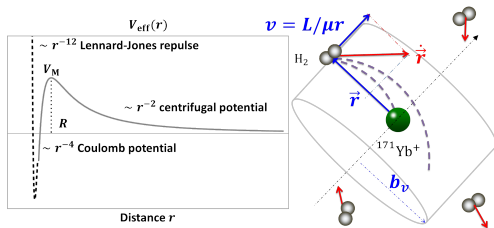




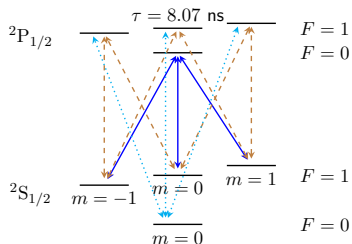
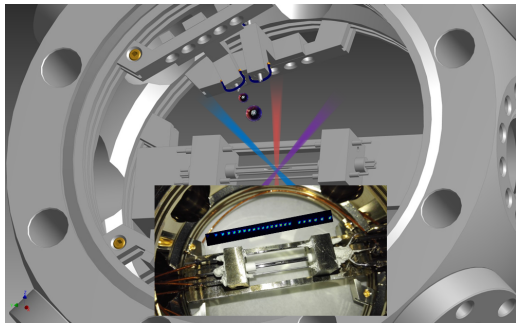
# System Construction



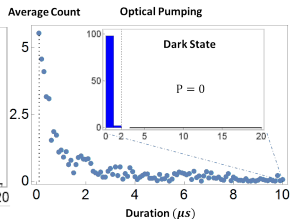
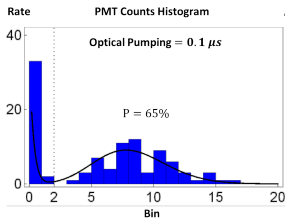
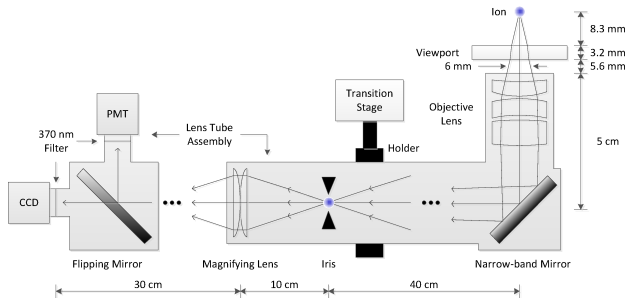
# Collision Estimation and UHV Preparation



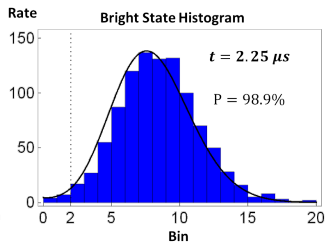
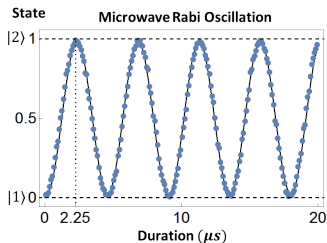
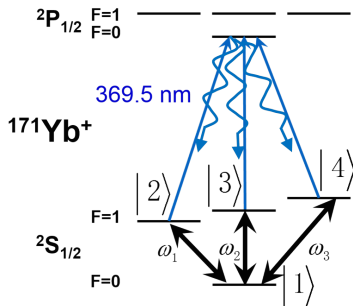
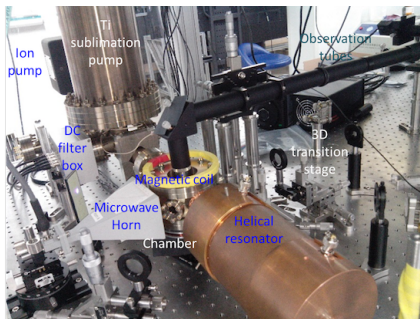
# Ionization and Doppler Cooling



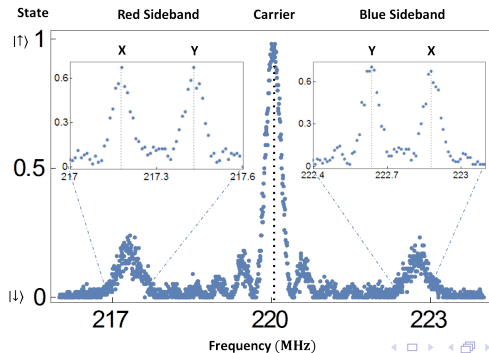
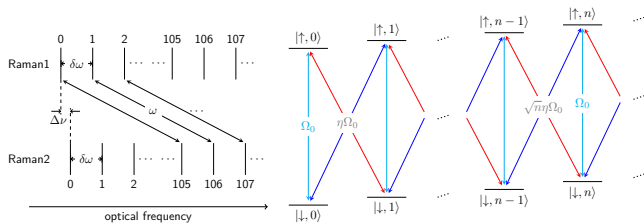
# State Detection and Initialization



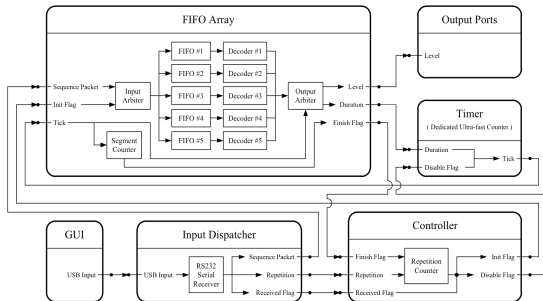
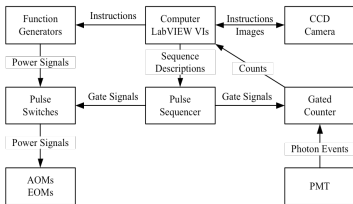
# Microwave Manipulation



# Pulse Laser Raman Transition



# Control System: Hardware



# Control System: Software

**Sequencer**

2013/11/25 00:58:15

State: **96.30**

Wavelength Calibration

- 397
- 866

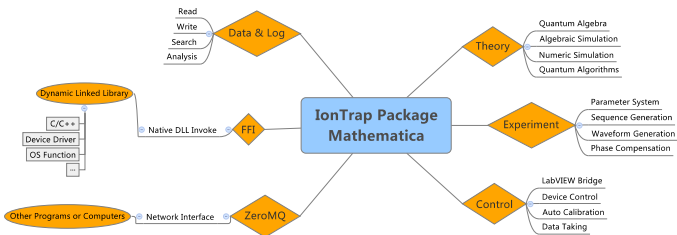
Sideband Cooling

RF = 4.0dbm

- Carrier1
- Repump
- RedZ1
- Carrier
- RedZ1
- RedZ
- BlueZ

Experiments

- Phonon Tomography





# Outline

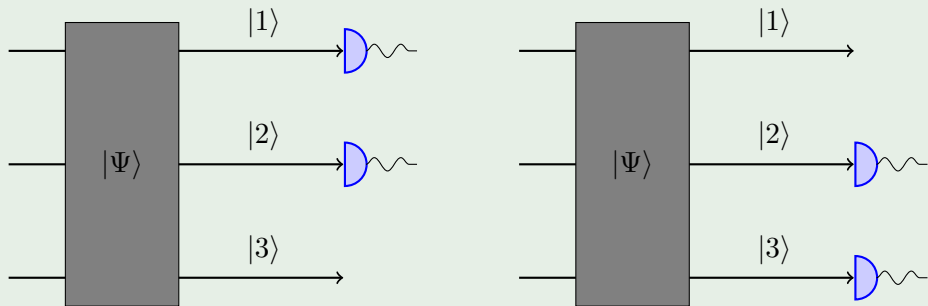
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# Non-Contextuality

## Definition

Observables' probability distribution are **independent of measurement**.

## Example



# Pentagram Inequality

## Hidden Variable Theory

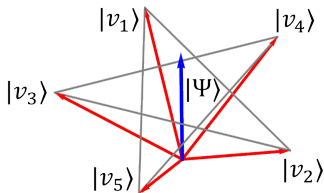
Let  $A_i$  be observables taking values  $\pm 1$ ,  $\langle \cdot \rangle$  denotes average value. Then

$$\langle A_1 A_2 \rangle + \langle A_2 A_3 \rangle + \langle A_3 A_4 \rangle + \langle A_4 A_5 \rangle + \langle A_5 A_1 \rangle \geq -3.$$

## Quantum Mechanics

Let  $A_i = I - 2 |v_i\rangle \langle v_i|$  be observables on state  $|\Psi\rangle$ . Then  $v(A_i) = \pm 1$ , and

$$\langle A_1 A_2 \rangle + \langle A_2 A_3 \rangle + \langle A_3 A_4 \rangle + \langle A_4 A_5 \rangle + \langle A_5 A_1 \rangle = 5 - 4\sqrt{5} \approx -3.944.$$



# Experimental Demonstration

$d = 3$  system is the most fundamental system shows contextuality.

## Previous Work

- $d \geq 4$ . Nature **460**, 494 (2009).
- $d = 3$ . Nature **474**, 490 (2011), PRL **109**,150401 (2012).

## Recently Experimental Demonstration

- $d = 3$ . PRL **110**, 070401 (2013)<sup>a</sup>.
- state-independent Kochen-Specker inequality
- with a single trapped ion (indivisible system, no entanglement)
- close detection efficiency loophole

<sup>a</sup>X. Zhang, et al., Phys. Rev. Lett. 110:070401 (2013)

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# State-independent Inequality<sup>2</sup>

## Hidden Variable Theory


Let  $A_i$  ( $i = 1, \dots, 13$ ) be observables taking values  $\pm 1$ . Then

$$\langle \chi_{13} \rangle := \sum_{i \in V} \mu_i \langle A_i \rangle - \sum_{(i,j) \in E} \mu_{ij} \langle A_i A_j \rangle - \sum_{(i,j,k) \in C} \mu_{ijk} \langle A_i A_j A_k \rangle \leq 25.$$

## Quantum Mechanics

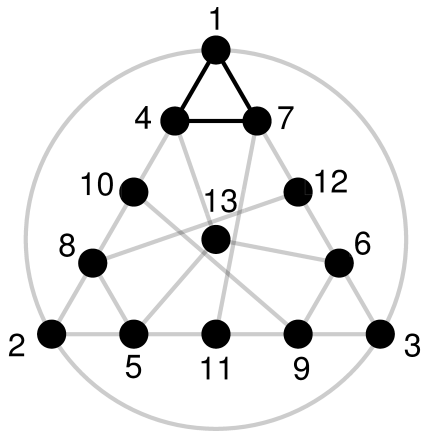
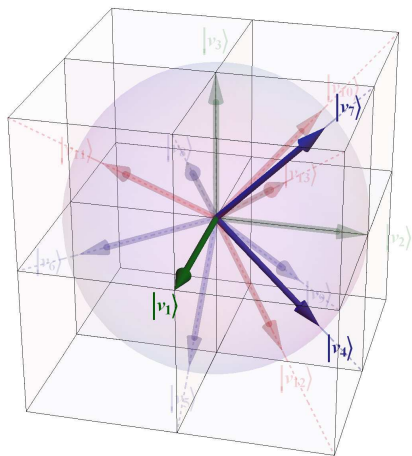
Let  $|v_i\rangle$  be basis vectors,  $A_i = I - 2|v_i\rangle\langle v_i|$  be observables. Then for any initial state  $|\Psi\rangle$ ,

$$\langle \chi_{13} \rangle = \frac{83}{3} \approx 27.67.$$

<sup>2</sup>A. Cabello, et al., Phys. Rev. Lett. 109:250402 (2012) 

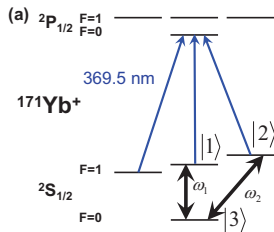


# Observables and Compatibility Relations





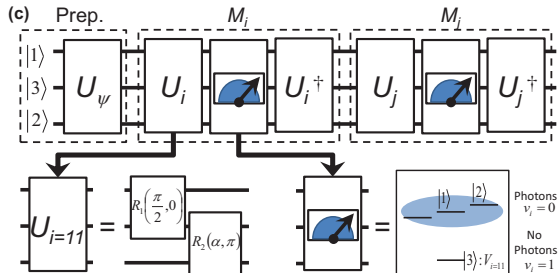
# Measurement Scheme



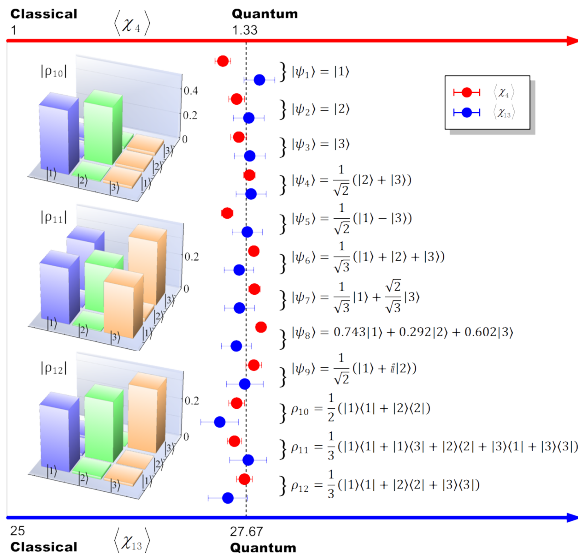
(b) Rotation matrices  $R_1(\theta_1, \phi_1)$  and  $R_2(\theta_2, \phi_2)$ :

$$R_1(\theta_1, \phi_1) = \begin{pmatrix} \cos(\theta_1/2) & 0 & e^{i\phi_1} \sin(\theta_1/2) \\ 0 & 1 & 0 \\ -e^{-i\phi_1} \sin(\theta_1/2) & 0 & \cos(\theta_1/2) \end{pmatrix}$$

$$R_2(\theta_2, \phi_2) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta_2/2) & -e^{-i\phi_2} \sin(\theta_2/2) \\ 0 & e^{i\phi_2} \sin(\theta_2/2) & \cos(\theta_2/2) \end{pmatrix}$$



# The Experimental Violation



$$\langle \chi_{13} \rangle = 27.38 \pm 0.21$$

# Summary

Quantum Contextuality is rooted in the fundamental structure of QM.

## Observed Experimental Violation

- $d = 3$ . PRL **110**, 070401 (2013).
- state-independent Kochen-Specker inequality
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- close detection efficiency loophole

## Outlook

- Application: “true” random number generator<sup>a</sup>
- Loophole-free: simultaneously measurement

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<sup>a</sup>U. Mark, X. Zhang, et al., Scientific Reports, 3:1627 (2013)

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# Majorana Particle<sup>3</sup>

- Majorana particle is its own antiparticle
- Whether *neutrinos* are Dirac or Majorana particles still remains open

## Majorana equation

$$i\hbar\gamma^\mu\partial_\mu\psi = mc\psi_c$$

$\gamma^\mu$  are Dirac matrices,  $\psi_c$  is the charge conjugate of the spinor  $\psi$ .

- Relativistic wave equation for fermions derived from first principles
- Preserves helicity and has no stationary solutions
- Relativistic quantum effects such as *Zitterbewegung*
- Time reversal and charge conjugation symmetries

<sup>3</sup>J. Casanova, et al., Phys. Rev. X, 1:021018 (2011)

## “Unphysical” Mapping

Majorana equation for (1 + 1) dimensions

$$i\hbar\partial_t\psi = c\hat{\sigma}_x\hat{p}_x\psi - mc^2\hat{\sigma}_y\psi^*$$

with “unphysical” operation mapping to enlarged space

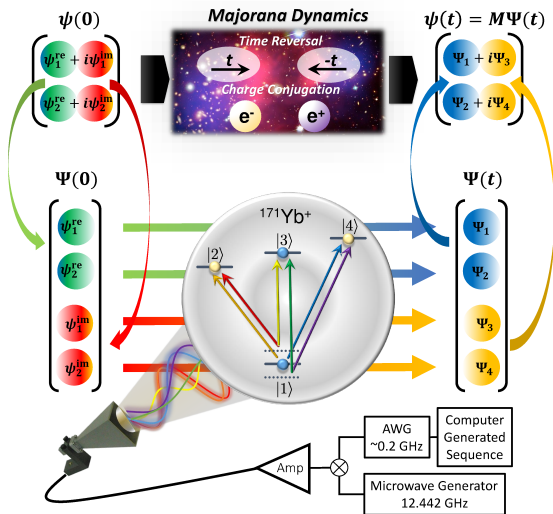
$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} \in \mathbb{C}_2 \rightarrow \Psi = \begin{pmatrix} \psi_1^r \\ \psi_2^r \\ \psi_1^i \\ \psi_2^i \end{pmatrix} \in \mathbb{R}_4$$

$$\psi = M\Psi = \begin{pmatrix} 1 & 0 & i & 0 \\ 0 & 1 & 0 & i \end{pmatrix} \Psi$$

becomes a (3 + 1)-dimensional Dirac equation

$$i\hbar\partial_t\Psi = [\hat{p}_xc(\mathbf{1} \otimes \sigma_x) - mc^2(\sigma_x \otimes \sigma_y)]\Psi$$

# Embedding Quantum Simulator<sup>4</sup>

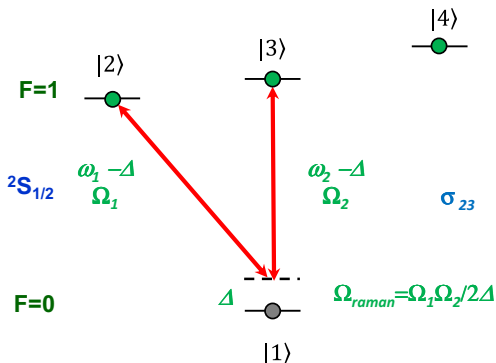


<sup>4</sup>X. Zhang, et al., Nature Communications, 6:7917 (2015)



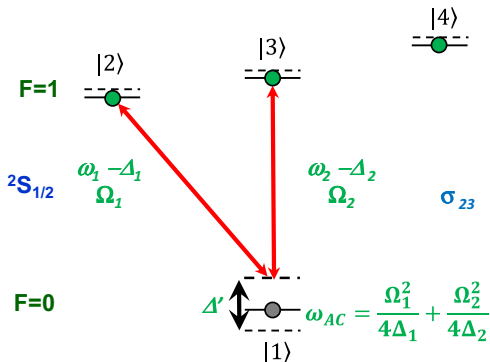
# Microwave Raman Transition

$$H_{\text{Majorana}} = \hat{p}_x c (\mathbf{1} \otimes \sigma_x) - mc^2 (\sigma_x \otimes \sigma_y) \rightarrow \underbrace{pc(\sigma_{12}^x + \sigma_{34}^x)}_{H_1} + \underbrace{mc^2(\sigma_{23}^y - \sigma_{14}^y)}_{H_2}$$



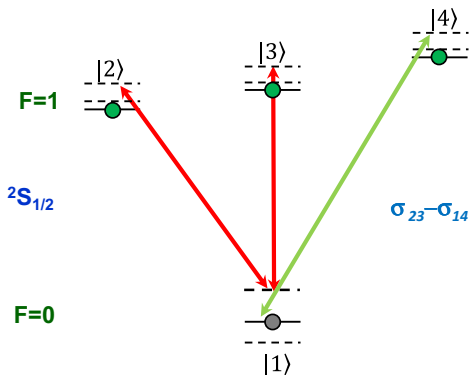
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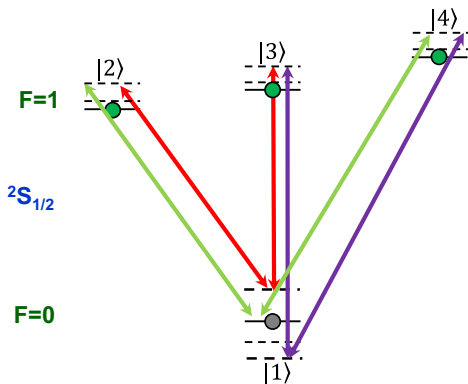
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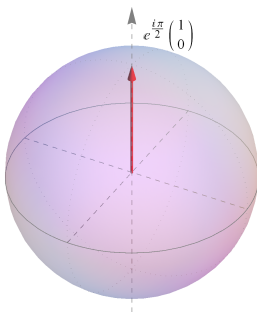


# Global Phase Effect

For parallel initial states with different global phase

$$|\psi_\theta(t=0)\rangle := e^{i\theta} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes |p\rangle$$

the fidelity defined as  $F(t) = |\langle \psi_\theta(t) | \psi_0(t) \rangle|^2$  is not conserved.

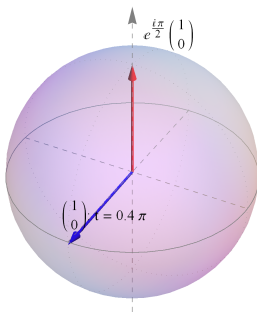


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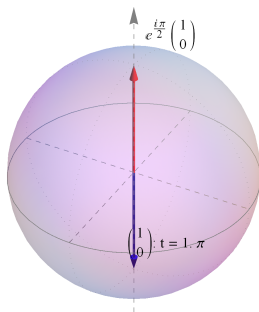


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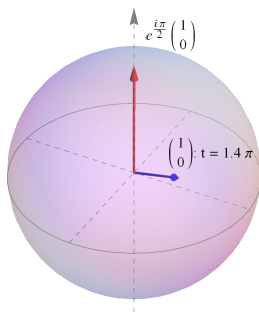


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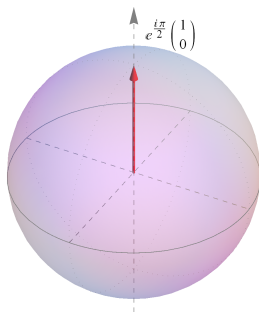


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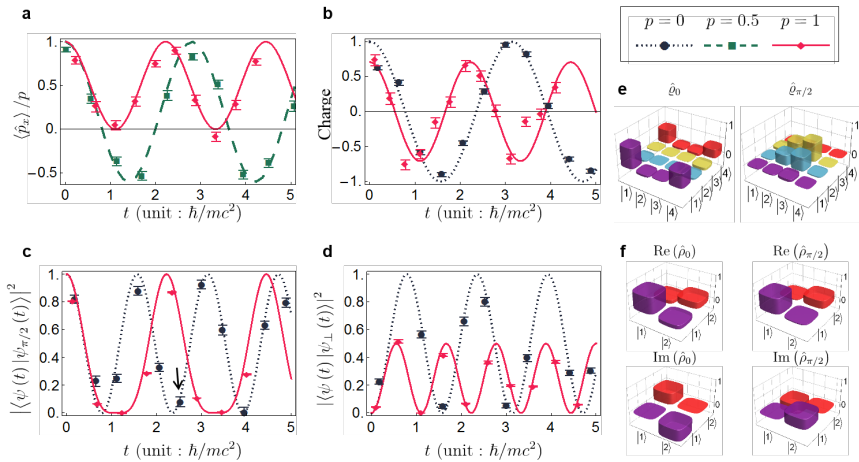
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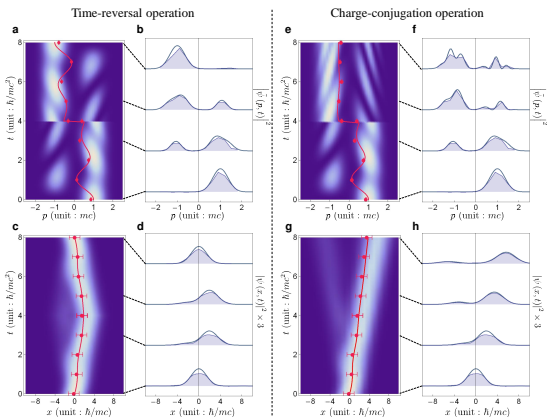
# Non-unitary Dynamics



# Symmetry Operations

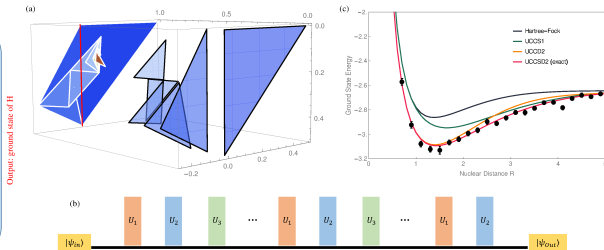
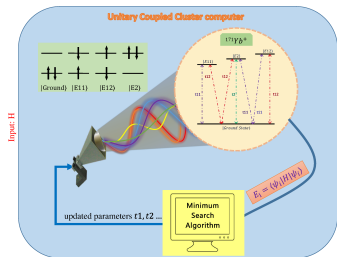
Apply symmetry operations at midpoint, with initial wave packet

$$\psi(x, t = 0) = (4\pi)^{-1/4} e^{-x^2/8} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} e^{ip_0 x/\hbar}$$



# Technique Application: Quantum Chemistry<sup>5</sup>

Ground energy of  $\text{HeH}^+$  calculated by Quantum Unitary Coupled Cluster



<sup>5</sup>Y. C. Shen, X. Zhang, et al., Phys. Rev. A. 95:020501 (2017)

# Summary

Realization of non-unitary dynamics and symmetry operations in a trapped-ion quantum simulator.

## Observed dynamics

- Global phase effect
- Orthogonality non-preservation
- Momentum *Zitterbewegung*
- Time reversal and charge conjugation

## Outlook

- Test discrete symmetry
- Anti-unitary operations with real momentum operator

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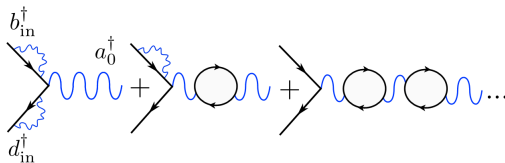
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# Simplified QFT Model<sup>6</sup>



## (1+1) QFT model

- Scalar fermions and bosons
- Fermion and anti-fermions interacting through bosonic field modes

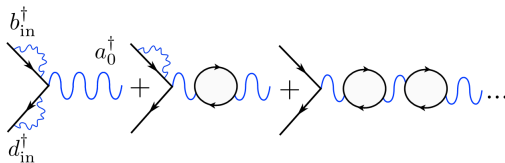
## Key features

- Fermion self-interaction process
- Particle creation and annihilation
- Non-perturbative regimes beyond Feynman diagrams

<sup>6</sup>J. Casanova, et al., Phys. Rev. Lett., 107, 260501 (2011)



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# Interaction Hamiltonian

$$\begin{aligned}
 H &= g \int dx \psi^\dagger(0, x) \psi(0, x) A(0, x) \\
 &\doteq g(t) (e^{i\delta t} b_{in}^\dagger d_{in}^\dagger a_0 + e^{-i(2\omega_0 + \delta)t} d_{in} b_{in} a_0) \\
 &\quad + g_1 e^{-i\omega_0 t} (b_{in}^\dagger b_{in} a_0 + d_{in} d_{in}^\dagger a_0) + \text{H.c.}
 \end{aligned}$$

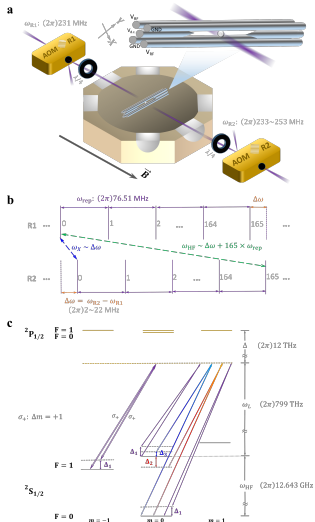
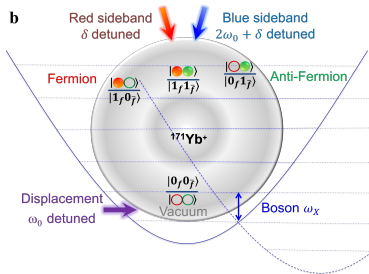
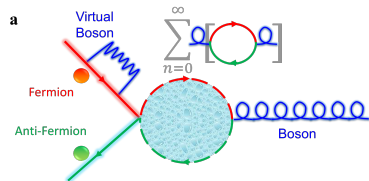
where  $\delta = \omega_f + \omega_{\bar{f}} - \omega_0$  and interaction strength  $g(t) = g_2 e^{-(t-T/2)^2/(2\sigma_t^2)}$ .

## Jordan-Wigner mapping

$$\begin{aligned}
 b_{in}^\dagger &= I \otimes \sigma^+, b_{in} = I \otimes \sigma^-, d_{in}^\dagger = \sigma^+ \otimes \sigma_z, d_{in} = \sigma^- \otimes \sigma_z \\
 H_I &= g_1 (|0_f 0_{\bar{f}}\rangle \langle 0_f 0_{\bar{f}}| + 2 |1_f 0_{\bar{f}}\rangle \langle 1_f 0_{\bar{f}}| + |1_f 1_{\bar{f}}\rangle \langle 1_f 1_{\bar{f}}|) \hat{a}_0 e^{-i\omega_0 t} \\
 &\quad - g(t) (|0_f 0_{\bar{f}}\rangle \langle 1_f 1_{\bar{f}}| \hat{a}_0^\dagger e^{-i\delta t} + |0_f 0_{\bar{f}}\rangle \langle 1_f 1_{\bar{f}}| \hat{a}_0 e^{-i(2\omega_0 + \delta)t}) + \text{H.c.}
 \end{aligned}$$

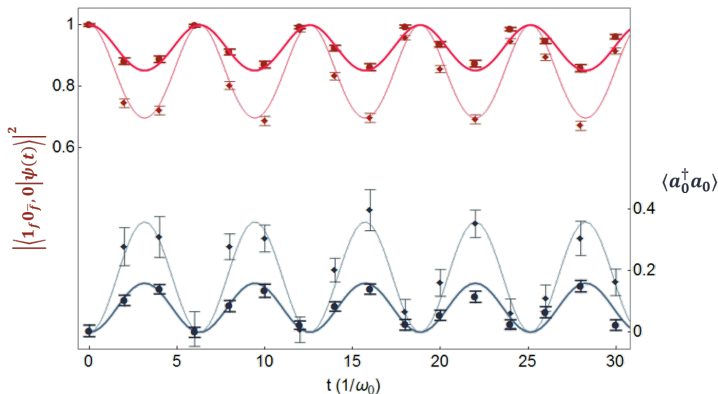
# Experimental Diagram

7



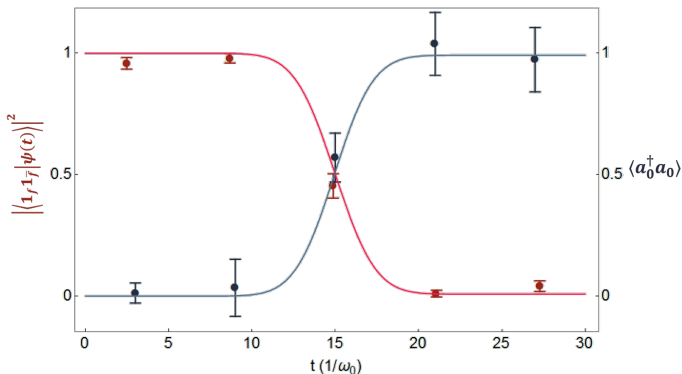
# Fermion Self-interaction

- Experiment parameter:  $g_1 = 0.15\omega_0, g_2 = 0, \sigma_t = 3/\omega_0$
- Initial state  $|1_f 0_{\bar{f}}, 0\rangle$ : one fermion state with no bosons
- Self-interaction dynamics:  $|1_f 0_{\bar{f}}, n\rangle \leftrightarrow |1_f 0_{\bar{f}}, n \pm 1\rangle$



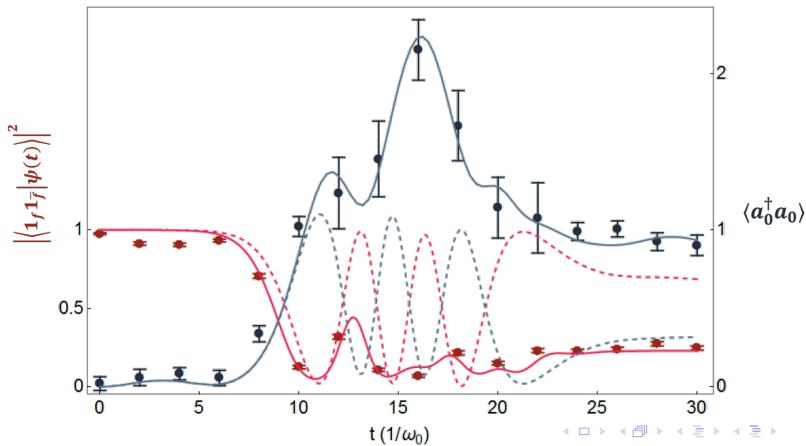
# Creation and Annihilation

- Experiment parameter:  $g_1 = 0.01\omega_0, g_2 = 0.21\omega_0, \sigma_t = 3/\omega_0$
- Initial state  $|1_f 1_{\bar{f}}, 0\rangle$ : fermion and antifermion state with no bosons
- Creation and annihilation dynamics:  $|1_f 1_{\bar{f}}, 0\rangle \leftrightarrow |0_f 0_{\bar{f}}, 1\rangle$



# Non-perturbative Regimes

- Experiment parameter:  $g_1 = 0.1\omega_0, g_2 = \omega_0, \sigma_t = 4/\omega_0$
- Initial state  $|1_f 1_{\bar{f}}, 0\rangle$ : fermion and antifermion state with no bosons
- Strong interaction coupling  $g_2 \geq \omega_0$
- Non-perturbative dynamics can't be calculated with Feynman diagram



# Summary

The first simulation of quantum field theory model with a trapped-ion quantum simulator.

## Observed dynamics

- Fermion self-interaction process
- Particle creation and annihilation
- Non-perturbative regimes beyond Feynman diagrams

## Outlook

- Extension to many field modes with ion chains
- Open quantum system Markov process
- 10 ions and 5 phonons/ion with dimension of  $2^{33} > 32\text{bit PC}$

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